

Comments on *Section 6.1.3 Supercritical  
Andronov-Hopf* of Izhikevich (2007)

Joaquín Rapela\*

October 3, 2016

# Contents

6	1	Introduction	1
7	2	Sketch of the derivation of the <i>approximate</i> topological normal	
8		form in Equations 6.8 and 6.9 of Izhikevich (2007)	2
9	3	Sketch of derivation of formulas for parameters $a$ and $d$ in Equa-	
10		tions 6.8 and 6.9 of Izhikevich (2007)	3
11	3.1	Rationale for Equations 6.5 and 6.6 in Izhikevich (2007) . . . . .	4
12	3.2	Solution to Problem 3.16 in Kuznetsov (2004) . . . . .	6
13	4	Derivation of formulas for $f(x, y)$ and $g(x, y)$ in Equation 6.6 of	
14		Izhikevich (2007)	9
15	5	Partial derivatives of $f(x, y)$ and $g(x, y)$ for the calculation of pa-	
16		rameters $a$ and $d$ in Equations 6.8 and 6.9 in Izhikevich (2007)	10
17	6	Parameters $a$ and $d$ in the example on page 172 of Izhikevich	
18		(2007)	11
19	7	Evidence for typo in the example on page 172 of Izhikevich	
20		(2007)	13

## 1 Introduction

The book Izhikevich (2007) is excellent. However, when I reached *Section 6.1.3 Supercritical Andronov-Hopf* I did not feel comfortable with the presentation of a few formulas without derivation or explanation of where these formulas came from or what do they meant. Specifically, I could not understand the rationale

---

\*rapela@ucsd.edu

for the change of variables in Equation 6.5, where the formula for parameter  $a$  in Equation 6.7 came from or what did it mean, how to derive the topological normal form in Equations 6.8 and 6.9, and how to compute parameter  $d$  in Equation 6.9.

Here I sketch the derivation of the topological normal form in Equations 6.8 and 6.9 and that of the formulas to calculate parameters  $a$  and  $d$  in these equations. In deriving the topological normal in Equations 6.8 and 6.9, I found that these equations do not represent a true topological normal form of a system undergoing an Andronov-Hopf bifurcation, but only an *approximation* of a true topological normal form. The formula for parameter  $a$  is given, without proof in Equation 6.7. On page 172 the author refers the reader to Exercise 6.17 to obtain the values of parameters  $a$  and  $d$  for an example  $I_{N_{a,p}} + I_K$  model. However, this exercise gives neither the formula nor the value of parameter  $d$ . Furthermore, applying the formulas we derive below for parameters  $a$  and  $d$  to the example  $I_{N_{a,p}} + I_K$  model on page 172 reveals a typo; the values of parameters  $a$  and  $d$  are interchanged in the penultimate paragraph on page 172, in the subsequent formula of the topological normal form of the  $I_{N_{a,p}} + I_K$  model, and the center and right panels of Figure 6.12 were computed with interchanged parameters.

Section 2 provides a sketch of the derivation of the *approximate* topological normal form in Equations 6.8 and 6.9. Section 3 sketches the derivation of the formulas of parameters  $a$  and  $d$  of the topological normal form for a system undergoing an Andronov-Hopf bifurcation. Section 4 derives the formulas for  $f(x, y)$  and  $g(x, y)$  in Equation 6.6. Section 6 calculates the values of parameters  $a$  and  $d$  for the example on page 172. Finally, Section 7 presents evidence for the typo in the example on page 172.

## 2 Sketch of the derivation of the *approximate* topological normal form in Equations 6.8 and 6.9 of Izhikevich (2007)

Consider a planar system

$$\dot{\mathbf{x}} = f(\mathbf{x}, \alpha), \quad \mathbf{x} \in \mathbb{R}^2, \quad \alpha \in \mathbb{R} \quad (1)$$

with a smooth function  $f$ , and an equilibrium at  $\mathbf{x} = 0$  for parameter  $\alpha = 0$ . This system undergoes an Andronov-Hopf bifurcation iff its eigenvalues are  $\lambda_{1,2} = \pm jw_0, w_0 > 0$ . From Section 3.5 Generic Hopf Bifurcation of Kuznetsov (2004) a planar system undergoing an Andronov-Hopf bifurcation can be rewritten as

$$\dot{\mathbf{x}} = A(\alpha)\mathbf{x} + F(\mathbf{x}, \alpha) \quad (2)$$

where  $F$  is a smooth vector function whose components  $F_{1,2}$  have Taylor expansions in  $\mathbf{x}$  starting with at least quadratic terms  $F = \mathcal{O}(\|\mathbf{x}\|^2)$ . Then, by

Lemma 3.3 in Kuznetsov (2004), the system in Equation 2 can be rewritten for sufficiently small  $|\alpha|$  a single complex equation

$$\dot{z} = \lambda(\alpha)z + g(z, \bar{z}, \alpha), \quad (3)$$

where  $g = \mathcal{O}(|z|^2)$  is a smooth function of  $(z, \bar{z}, \alpha)$ . Now from Lemma 3.6 (Poincaré normal form for the Hopf bifurcation) in Kuznetsov (2004) Equation 3 can be transformed, by an invertible parameter-dependent change of complex coordinate, for sufficiently small  $|\alpha|$ , into

$$\dot{w} = \lambda(\alpha)w + c_1(\alpha)w^2\bar{w} + \mathcal{O}(|w|^4). \quad (4)$$

Denoting  $\lambda(\alpha) = c(\alpha) + jd(\alpha)$  and  $c_1(\alpha) = a(\alpha) + jd(\alpha)$ , by Problem 6.4 in Izhikevich (2007), the complex-valued Equation 4 can be written in polar coordinates as

$$\begin{aligned} \dot{r} &= c(\alpha)r + a(\alpha)r^3, \\ \dot{\phi} &= w(\alpha) + d(\alpha)r^2 \end{aligned} \quad (5)$$

The system in Equation 5 is a true topological normal form of the system in Equation 1 undergoing an Andronov-Hopf bifurcation, because the former system can be obtained from the latter system by a change of variables that is an homeomorphism. However, the topological normal form in Equations 6.8 and 6.9 of Izhikevich (2007) is an approximation of the true topological normal form using  $a(0)$  and  $d(0)$  instead of  $a(\alpha)$  and  $d(\alpha)$  in Equation 5.

To obtain exact values of  $a$  and  $d$  independent of any parameter, one could use a topological normal form as in Equation 4 but with a coefficient of the cubic term independent of any parameter (e.g.,  $c_1(\alpha) = c_1$  in Equation 4). This new topological normal form is given in Lemma 3.7 of Kuznetsov (2004). In this lemma Kuznetsov (2004) attains a topological normal form where the coefficient of the cubic term is  $\text{sign}(l_1(\beta))$ . Because  $l_1(\beta)$  is a continuous function of  $\beta$ , with  $l_1(0) \neq 0$ , for small values of  $\beta$  the sign of  $l_1(\beta)$  is constant (i.e.,  $\text{sign}(l_1(\beta)) = s$ ). Therefore, by replacing  $\text{sign}(l_1(\beta))$  with  $s$  in the previous topological normal form, for small values of  $\beta$ , Kuznetsov attains a new topological normal form with the coefficient of the cubic term independent of any parameter.

### 3 Sketch of derivation of formulas for parameters $a$ and $d$ in Equations 6.8 and 6.9 of Izhikevich (2007)

As shown by the end of Section 2 parameters  $a$  and  $d$  in Equations 6.8 and 6.9 of Izhikevich (2007) are  $a = \text{Re}\{c_1(0)\}$  and  $d = \text{Im}\{c_1(0)\}$ , with  $c_1(0)$  given in

Equation 4. To calculate  $c_1(0)$  we first apply the change of variables in Equation 6.5 of Izhikevich (2007) to convert the system undergoing an Andronov-Hopf bifurcation in Equation 6.4 of Izhikevich (2007), with parameter  $b = 0$ , into the form in Equation 6.6 of Izhikevich (2007). Then we solve Problem (3.16) in Kuznetsov (2004) to obtain  $c_1(0)$  from the system in Equation 6.6 of Izhikevich (2007).

Section 3.1 provides the rationale for Equations 6.5 and 6.6 of Izhikevich (2007). Section 3.2 sketches the solution of Problem 3.16 in Kuznetsov (2004).

### 3.1 Rationale for Equations 6.5 and 6.6 in Izhikevich (2007)

The rationale for the change of variables in Equation 6.5 of Izhikevich (2007) is given in Lemma 3.3 of Kuznetsov (2004). This lemma proves that a system of differential equations

$$\dot{\mathbf{x}} = A(\alpha)\mathbf{x} + F(\mathbf{x}, \alpha)$$

where  $F : \mathbb{R}^2 \times \mathbb{R} \rightarrow \mathbb{R}^2$  is a smooth vector function whose components have Taylor expansions in  $\mathbf{x}$  starting with at least quadratic terms, can be written as a single differential equation on a complex variable  $z$ :

$$\dot{z} = \lambda(\alpha)z + g(z, \bar{z}, \alpha) : \quad (6)$$

with  $\lambda(\alpha)$  being the eigenvalue of  $A(\alpha)$  corresponding to eigenvector  $\mathbf{q}(\alpha)$  (i.e.,  $A(\alpha)\mathbf{q}(\alpha) = \lambda(\alpha)\mathbf{q}(\alpha)$ ) and  $z$  being the coordinate associated with  $\mathbf{q}(\alpha)$  in the expansion of  $\mathbf{x}$  in the base  $\{\mathbf{q}(\alpha), \bar{\mathbf{q}}(\alpha)\}$ ,

$$\mathbf{x} = z\mathbf{q}(\alpha) + \bar{z}\bar{\mathbf{q}}(\alpha), \quad (7)$$

and

$$g(z, \bar{z}, \alpha) = \langle \mathbf{p}(\alpha), F(z\mathbf{q}(\alpha) + \bar{z}\bar{\mathbf{q}}(\alpha), \alpha) \rangle \quad (8)$$

with  $\{\mathbf{p}(\alpha), \bar{\mathbf{p}}(\alpha)\}$  a base of eigenvectors of  $A(\alpha)^T$  biorthogonal with the base  $\{\mathbf{q}(\alpha), \bar{\mathbf{q}}(\alpha)\}$  of  $A(\alpha)$ . The Taylor series expansion of the right-hand-side of the differential equations in Equation 6.4 of Izhikevich (2007), for  $b = 0$  can be written as:

$$\begin{bmatrix} \dot{v} \\ \dot{u} \end{bmatrix} = A(0) \begin{bmatrix} v \\ u \end{bmatrix} + \begin{bmatrix} \tilde{F}(v, u, 0) \\ \tilde{G}(v, u, 0) \end{bmatrix} \quad (9)$$

with  $A(0)$  being the jacobian matrix of the right hand side of the differential equations in Equation 6.4 of Izhikevich (2007)

$$A(0) = \begin{bmatrix} F_u & F_v \\ G_u & G_v \end{bmatrix} \quad (10)$$

117 where  $F_u, F_v, G_u, G_v$  are partial derivatives evaluated at  $(u, v, b) = (0, 0, 0)$ .  
 118 Equation 9 holds because it is assumed that  $(v, u) = (0, 0)$  is an equilibrium  
 119 for  $b = 0$  of the system of differential equations in Equation 6.4 of Izhikevich  
 120 (2007), then  $F(0, 0, 0) = G(0, 0, 0) = 0$ , and the Taylor series expansion of the  
 121 right-hand-side of Equation 6.4 in Izhikevich (2007) has no constant term.

122 Now, because it is assumed that the system in Equation 6.4 in Izhikevich  
 123 (2007) undergoes an Andronov-Hopf bifurcation at the equilibrium  $(v, u) =$   
 124  $(0, 0)$ , by the non-hyperbolicity condition of the Andronov-Hopf bifurcation,  
 125 the eigenvalues of  $A(0)$  in Equation 9 are  $\pm jw$ . Then by Equation 6:

$$\dot{z} = jwz + g(z, \bar{z}, 0) \quad (11)$$

126 with  $z$  being the coordinate associated with  $\mathbf{q}(0)$  in the expansion of  $[v, u]^T$  in  
 127 the base  $\{\mathbf{q}(0), \bar{\mathbf{q}}(0)\}$ . That is, from Equation 7,

$$\begin{bmatrix} v \\ u \end{bmatrix} = z\mathbf{q}(0) + \bar{z}\bar{\mathbf{q}}(0) \quad (12)$$

128 Separating the real and imaginary parts of  $z$  in Equation 11 we obtain:

$$\begin{aligned} \text{Re}\{\dot{z}\} &= -w\text{Im}\{z\} + \text{Re}\{g(z, \bar{z}, 0)\} \\ \text{Im}\{\dot{z}\} &= w\text{Re}\{z\} + \text{Im}\{g(z, \bar{z}, 0)\} \end{aligned}$$

129 and calling

$$x = \text{Re}\{z\}, \quad (13)$$

$$y = \text{Im}\{z\}, \quad (14)$$

$$f(x, y) = \text{Re}\{g(z, \bar{z}, 0)\}, \quad (15)$$

$$g(x, y) = \text{Im}\{g(z, \bar{z}, 0)\}, \quad (16)$$

130 we obtain

$$\begin{aligned} \dot{x} &= -wy + f(x, y) \\ \dot{y} &= wx + g(x, y) \end{aligned}$$

131 as in Equation 6.6 in Izhikevich (2007).

132 To express  $x$  and  $y$  as a function of  $v$  and  $u$ , as in Equation 6.5 of Izhikevich  
 133 (2007), from Equations 13 and 14, we want to express  $z$  as a function of  $v$  and

134 *u*. One can build the biorthogonal base  $\{\mathbf{p}(0), \bar{\mathbf{p}}(0)\}$  corresponding to the base  
 135 of eigenvectors of  $A(0)$ ,  $\{\mathbf{q}(0), \bar{\mathbf{q}}(0)\}$  and, from Equation 12, compute  $z$  as

$$z = \langle \mathbf{p}(0), \begin{bmatrix} v \\ u \end{bmatrix} \rangle \quad (17)$$

136 Because  $\bar{\mathbf{q}}(0)$  is the eigenvector of  $A(0)$  (Equation 10) with eigenvalue  $-jw$ ,  
 137 we have

$$\begin{bmatrix} F_v + jw & F_u \\ G_v & G_u + jw \end{bmatrix} \begin{bmatrix} \bar{q}_1(0) \\ \bar{q}_2(0) \end{bmatrix} = 0$$

138 Then

$$(F_v + jw)\bar{q}_1(0) + F_u\bar{q}_2(0) = 0 \quad (18)$$

139 Due to the biorthogonality of the bases  $\{\mathbf{p}(0), \bar{\mathbf{p}}(0)\}$  and  $\{\mathbf{q}(0), \bar{\mathbf{q}}(0)\}$  we have  
 140  $\langle \mathbf{p}(0), \bar{\mathbf{q}}(0) \rangle = 0$ . Thus, from Equation 18, we obtain

$$\mathbf{p}(0) = k \begin{bmatrix} F_v - jw \\ F_u \end{bmatrix} \quad (19)$$

141 where  $k$  is a complex constant. Now, from Equations 17 and 19

$$\begin{aligned} z &= \langle \mathbf{p}(0), \begin{bmatrix} v \\ u \end{bmatrix} \rangle \\ &= \bar{k}v(F_v + jw) + \bar{k}uF_u \\ &= \bar{k}(uF_u + vF_v) + j\bar{k}wv \end{aligned} \quad (20)$$

142 By using  $k = j/w$  in Equation 20 we obtain  $z = v - j(uF_u + vF_v)/w$ . Then  
 143  $x = \text{Re}\{z\} = v$ , and  $y = \text{Im}\{z\} = -(uF_u + vF_v)/w = -(uF_u + xF_v)/w$ , or  
 144  $v = x$  and  $F_u u = -F_v x - wy$  as in Equation 5.5 of Izhikevich (2007).

### 145 **3.2 Solution to Problem 3.16 in Kuznetsov (2004)**

146 Given a system as that in Equation 6.6 of Izhikevich (2007), Problem 3.16 in  
 147 Kuznetsov (2004) asks to compute  $c_1(0)$ . The solution is  $c_1(0) = a(0) + jd(0) =$   
 148  $a + jd$  with

$$\begin{aligned} a &= \frac{1}{16}((f_{xxx} + g_{xxy} + f_{xyy} + g_{yyy}) + \\ &\quad \frac{1}{w}(g_{yy}f_{yy} - g_{xx}f_{xx} + f_{xy}(f_{xx} + f_{yy}) - g_{xy}(g_{xx} + g_{yy}))), \end{aligned} \quad (21)$$

$$\begin{aligned}
d = & \frac{1}{16}((g_{xxx} + g_{xyy} - f_{xxy} - f_{yyx}) + \\
& \frac{1}{3w}((f_{xy}g_{xx} + g_{xy}f_{yy}) - 2(g_{yy}^2 + f_{xx}^2) - 2(f_{xy}^2 + g_{xy}^2) - \\
& 5(g_{xx}^2 + f_{yy}^2) + 5(f_{xy}g_{yy} + g_{xy}f_{xx}) - \\
& 5(g_{xx}g_{yy} + f_{xx}f_{yy}))).
\end{aligned} \tag{22}$$

149 To derive this solution, we first rewrite Equation 6.6 of Izhikevich (2007) in  
150 matrix form as

$$\begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} = A(0) \begin{bmatrix} x \\ y \end{bmatrix} + F\left(\begin{bmatrix} x \\ y \end{bmatrix}, 0\right) \tag{23}$$

151 with

$$A_0 = \begin{bmatrix} 0 & -w \\ w & 0 \end{bmatrix}, \tag{24}$$

$$F\left(\begin{bmatrix} x \\ y \end{bmatrix}, 0\right) = \begin{bmatrix} f(x, y) \\ g(x, y) \end{bmatrix}, \tag{25}$$

152 and  $f$  and  $g$  having Taylor series expansions starting with at least quadratic  
153 terms

$$\begin{aligned}
f(x, y) = & \frac{1}{2}f_{xx}x^2 + f_{xy}xy + \frac{1}{2}f_{yy}y^2 + \\
& \frac{1}{6}f_{xxx}x^3 + \frac{1}{2}f_{xxy}x^2y + \frac{1}{2}f_{xyy}xy^2 + \frac{1}{6}f_{yyy}y^3 + \dots,
\end{aligned} \tag{26}$$

$$\begin{aligned}
g(x, y) = & \frac{1}{2}g_{xx}x^2 + g_{xy}xy + \frac{1}{2}g_{yy}y^2 + \\
& \frac{1}{6}g_{xxx}x^3 + \frac{1}{2}g_{xxy}x^2y + \frac{1}{2}g_{xyy}xy^2 + \frac{1}{6}g_{yyy}y^3 + \dots
\end{aligned} \tag{27}$$

154 with all partial derivatives of  $f$  and  $g$  evaluated at the origin. From Equa-  
155 tion 3.18 of Kuznetsov (2004),

$$c_1(0) = \frac{j}{2w} \left( g_{20}g_{11} - 2|g_{11}|^2 - \frac{1}{3}|g_{02}|^2 \right) + \frac{g_{21}}{2}, \tag{28}$$

156 where, from p. 91 of Kuznetsov (2004),  $g_{kl}$  are the coefficients of the Taylor  
157 series expansion of the function  $g(z, \bar{z}, 0)$  in Equation 8,

$$g(z, \bar{z}, 0) = \sum_{k+l \geq 2} \frac{1}{k!l!} g_{kl}(0) z^k \bar{z}^l. \tag{29}$$

158 To calculate  $c_1(0)$  we will expand  $g(z, \bar{z}, 0)$  in Taylor and extract the coeffi-  
 159 cients  $g_{11}$ ,  $g_{20}$ ,  $g_{02}$ , and  $g_{21}$  needed in Equation 28. We choose  $\mathbf{p}(0) = [-j, -1]^T$   
 160 and  $\mathbf{q}(0) = [-j/2, -1/2]^T$ . Then, from Equation 8,

$$g(z, \bar{z}, 0) = jf(z\mathbf{q}(0) + \bar{z}\bar{\mathbf{q}}(0)) - g(z\mathbf{q}(0) + \bar{z}\bar{\mathbf{q}}(0)). \quad (30)$$

161 Inserting

$$z\mathbf{q}(0) + \bar{z}\bar{\mathbf{q}}(0) = -\frac{1}{2} \begin{bmatrix} jz - j\bar{z}, \\ z + \bar{z} \end{bmatrix} = \begin{bmatrix} x \\ y \end{bmatrix}$$

162 into Equations 26 and 27 we obtain

$$\begin{aligned} f(z\mathbf{q}(0) + \bar{z}\bar{\mathbf{q}}(0)) = & \left(-\frac{1}{8}f_{xx} + \frac{1}{8}f_{yy} + j\frac{1}{4}f_{xy}\right)z^2 + \\ & \left(\frac{1}{4}f_{xx} + \frac{1}{4}f_{yy}\right)z\bar{z} + \\ & \left(-\frac{1}{8}f_{xx} + \frac{1}{8}f_{yy} - j\frac{1}{4}f_{xy}\right)\bar{z}^2 + \\ & \left(\frac{1}{16}f_{xxy} - \frac{1}{48}f_{yyy} + j\left(\frac{1}{48}f_{xxx} - \frac{1}{16}f_{xyy}\right)\right)z^3 + \\ & \left(-\frac{1}{16}f_{xxy} - \frac{1}{16}f_{yyy} - j\left(\frac{1}{16}f_{xxx} + \frac{1}{16}f_{xyy}\right)\right)z^2\bar{z} + \\ & \left(-\frac{1}{16}f_{xxy} - \frac{1}{16}f_{yyy} - j\left(\frac{1}{16}f_{xxx} + \frac{1}{16}f_{xyy}\right)\right)z\bar{z}^2 + \\ & \left(\frac{1}{16}f_{xxy} - \frac{1}{48}f_{yyy} - j\left(\frac{1}{48}f_{xxx} - \frac{1}{16}f_{xyy}\right)\right)\bar{z}^3 \end{aligned} \quad (31)$$

$$\begin{aligned} g(z\mathbf{q}(0) + \bar{z}\bar{\mathbf{q}}(0)) = & \left(-\frac{1}{8}g_{xx} + \frac{1}{8}g_{yy} + j\frac{1}{4}g_{xy}\right)z^2 + \\ & \left(\frac{1}{4}g_{xx} + \frac{1}{4}g_{yy}\right)z\bar{z} + \\ & \left(-\frac{1}{8}g_{xx} + \frac{1}{8}g_{yy} - j\frac{1}{4}g_{xy}\right)\bar{z}^2 + \\ & \left(\frac{1}{16}g_{xxy} - \frac{1}{48}g_{yyy} + j\left(\frac{1}{48}g_{xxx} - \frac{1}{16}g_{xyy}\right)\right)z^3 + \\ & \left(-\frac{1}{16}g_{xxy} - \frac{1}{16}g_{yyy} - j\left(\frac{1}{16}g_{xxx} + \frac{1}{16}g_{xyy}\right)\right)z^2\bar{z} + \\ & \left(-\frac{1}{16}g_{xxy} - \frac{1}{16}g_{yyy} - j\left(\frac{1}{16}g_{xxx} + \frac{1}{16}g_{xyy}\right)\right)z\bar{z}^2 + \\ & \left(\frac{1}{16}g_{xxy} - \frac{1}{48}g_{yyy} - j\left(\frac{1}{48}g_{xxx} - \frac{1}{16}g_{xyy}\right)\right)\bar{z}^3 \end{aligned} \quad (32)$$

163 Using Equations 31 and 32 in Equation 16 we obtain the Taylor series ex-  
 164 pansion of  $g(z, \bar{z}, 0)$ , and from this expansion we extract the coefficients  $g_{11}$ ,  $g_{20}$ ,  
 165  $g_{02}$ , and  $g_{21}$  needed in Equation 28. The obtained coefficients are



$$g_{20} = g_{z^2} = \frac{1}{4}((g_{xx} - g_{yy} - 2f_{xy}) + j(f_{yy} - f_{xx} - 2g_{xy})) \quad (33)$$

$$g_{11} = g_{z\bar{z}} = \frac{1}{4}((-g_{xx} - g_{yy}) + j(f_{yy} + f_{xx})) \quad (34)$$

$$g_{02} = g_{\bar{z}^2} = \frac{1}{4}((g_{xx} - g_{yy} + 2f_{xy}) + j(f_{yy} - f_{xx} + 2g_{xy})) \quad (35)$$

$$g_{21} = g_{z^2\bar{z}} = \frac{1}{8}((f_{xxx} + g_{xxy} + f_{xyy} + g_{yyy}) + j(-f_{xxy} + g_{xx} - f_{yy} + g_{xy})) \quad (36)$$

166 From Equation 16 and Equations 33-36 we obtain

$$\begin{aligned} a = \text{Re}\{c_1(0)\} &= -\frac{1}{2w} \text{Im}\{g_{20}g_{11}\} + \frac{1}{2} \text{Re}\{g_{21}\} = \\ &= \frac{1}{16}((f_{xxx} + g_{xxy} + f_{xyy} + g_{yyy}) + \\ &\quad \frac{1}{w}(g_{yy}f_{yy} - g_{xx}f_{xx} + f_{xy}(f_{xx} + f_{yy}) - g_{xy}(g_{xx} + g_{yy}))), \\ d = \text{Im}\{c_1(0)\} &= \frac{1}{2w}(\text{Re}\{g_{20}g_{11}\} - 2|g_{11}|^2 - \frac{1}{3}|g_{02}|^2) + \frac{1}{2} \text{Im}\{g_{21}\} = \\ &= \frac{1}{16}((g_{xxx} + g_{xxy} - f_{xxy} - f_{yyy}) + \\ &\quad \frac{1}{3w}((f_{xy}g_{xx} + g_{xy}f_{yy}) - 2(g_{yy}^2 + f_{xx}^2) - 2(f_{xy}^2 + g_{xy}^2) - \\ &\quad 5(g_{xx}^2 + f_{yy}^2) + 5(f_{xy}g_{yy} + g_{xy}f_{xx}) - \\ &\quad 5(g_{xx}g_{yy} + f_{xx}f_{yy}))), \end{aligned}$$

167 as in Equations 21 and 22.

## 168 **4 Derivation of formulas for $f(x, y)$ and $g(x, y)$ in** 169 **Equation 6.6 of Izhikevich (2007)**

170 Equations 15 and 16 give formulas for  $f(x, y)$  and  $g(x, y)$  in Equation 6.6 of  
171 Izhikevich (2007). However, to use these formulas one first needs to compute the  
172 function  $g(z, \bar{z}, 0)$ . Below we derive formulas for  $f(x, y)$  and  $g(x, y)$  as a function  
173 of the functions  $F(v, u)$  and  $G(v, u)$  in Equation 6.4 of Izhikevich (2007), that  
174 do not require function  $g(z, \bar{z}, 0)$ .

175 From the change of variables in Equation 6.5 of Izhikevich (2007) we have

$$\begin{aligned} \dot{x} = \dot{v} &= F(v, u, 0) = F(x, -\frac{1}{F_n}(F_v x + wy)) \\ &= -wy + F(x, -\frac{1}{F_n}(F_v x + wy), 0) + wy = -wy + f(x, y), \end{aligned}$$

176 with

$$f(x, y) = F(x, -\frac{1}{F_n}(F_v x + wy), 0) + wy. \quad (37)$$

177 Also, from the change of variables in Equation 6.5 of Izhikevich (2007) we  
178 obtain, for  $w \neq 0$ ,

$$y = -\frac{1}{w}(F_u u + F_v x).$$

179 Then

$$\begin{aligned} \dot{y} &= -\frac{1}{w}(F_u \dot{u} + F_v \dot{x}) \\ &= -\frac{1}{w}(F_u G(v, u, 0) + F_v F(v, u, 0)) \\ &= -\frac{1}{w}(F_u G(x, -\frac{1}{F_n}(F_v x + wy), 0) + F_v F(x, -\frac{1}{F_n}(F_v x + wy), 0)) \\ &= wx - \frac{1}{w}(F_u G(x, -\frac{1}{F_n}(F_v x + wy), 0) + F_v F(x, -\frac{1}{F_n}(F_v x + wy), 0)) - wx \\ &= wx + g(x, y), \end{aligned}$$

180 with

$$\begin{aligned} g(x, y) &= -\frac{1}{w}(F_u G(x, -\frac{1}{F_n}(F_v x + wy), 0) + F_v F(x, -\frac{1}{F_n}(F_v x + wy), 0)) \\ &\quad - wx. \end{aligned} \quad (38)$$

## 181 5 Partial derivatives of $f(x, y)$ and $g(x, y)$ for the 182 calculation of parameters $a$ and $d$ in Equa- 183 tions 6.8 and 6.9 in Izhikevich (2007)

184 Although tedious, it is not difficult to compute the partial derivatives of  $f(x, y)$   
185 and  $g(x, y)$  in Equation 37 and 38, respectively, and evaluate them at  $x = 0$  and  
186  $y = 0$ . We obtained:

$$\begin{aligned} f_{xx} &= F_{VV} + 2F_{Vn} \left(-\frac{F_V}{F_n}\right) + F_{nn} \left(-\frac{F_V}{F_n}\right)^2 \\ f_{xxx} &= F_{VVV} + 3F_{VVn} \left(-\frac{F_V}{F_n}\right) + 3F_{Vnn} \left(-\frac{F_V}{F_n}\right)^2 + F_{nnn} \left(-\frac{F_V}{F_n}\right)^3 \end{aligned}$$

$$\begin{aligned}
f_{xxy} &= \left(-\frac{w}{F_n}\right) \left(F_{VVn} + 2F_{Vnn} \left(-\frac{F_V}{F_n}\right) + F_{nnn} \left(-\frac{F_V}{F_n}\right)^2\right) \\
f_{yy} &= F_{nn} \left(-\frac{w}{F_n}\right)^2 \\
f_{yyy} &= F_{nnn} \left(-\frac{w}{F_n}\right)^3 \\
f_{xyy} &= \left(-\frac{w}{F_n}\right)^2 \left(F_{Vnn} + F_{nnn} \left(-\frac{F_V}{F_n}\right)\right) \\
f_{xy} &= \left(-\frac{w}{F_n}\right) \left(F_{Vn} + F_{nn} \left(-\frac{F_V}{F_n}\right)\right) \\
g_{xx} &= -\frac{1}{w} \left(F_V \left(F_{VV} + 2F_{Vn} \left(-\frac{F_V}{F_n}\right) + F_{nn} \left(-\frac{F_V}{F_n}\right)^2\right) + \right. \\
&\quad \left. F_n \left(G_{VV} + 2G_{Vn} \left(-\frac{F_V}{F_n}\right) + G_{nn} \left(-\frac{F_V}{F_n}\right)^2\right)\right) \\
g_{xxx} &= -\frac{1}{w} \left(F_V \left(F_{VVV} + 3F_{VVn} \left(-\frac{F_V}{F_n}\right) + 3F_{Vnn} \left(-\frac{F_V}{F_n}\right)^2 + F_{nnn} \left(-\frac{F_V}{F_n}\right)^3\right) + \right. \\
&\quad \left. F_n \left(G_{VVV} + 3G_{VVn} \left(-\frac{F_V}{F_n}\right) + 3G_{Vnn} \left(-\frac{F_V}{F_n}\right)^2 + G_{nnn} \left(-\frac{F_V}{F_n}\right)^3\right)\right) \\
g_{xxy} &= \frac{1}{F_n} \left(F_V \left(F_{VVn} + 2F_{Vnn} \left(-\frac{F_V}{F_n}\right) + F_{nnn} \left(-\frac{F_V}{F_n}\right)^2\right) + \right. \\
&\quad \left. F_n \left(G_{VVn} + 2G_{Vnn} \left(-\frac{F_V}{F_n}\right) + G_{nnn} \left(-\frac{F_V}{F_n}\right)^2\right)\right) \\
g_{yy} &= -\frac{w}{F_n^2} (F_V F_{nn} + F_n G_{nn}) \\
g_{yyy} &= \frac{w^2}{F_n^3} (F_V F_{nnn} + F_n G_{nnn}) \\
g_{xyy} &= -\frac{w}{F_n^2} (F_V (F_{Vnn} + F_{nnn}) + F_n (G_{Vnn} + G_{nnn})) \\
g_{xy} &= \frac{1}{F_n} \left(F_V \left(F_{Vn} + F_{nn} \left(-\frac{F_V}{F_n}\right)\right) + F_n \left(G_{Vn} + G_{nn} \left(-\frac{F_V}{F_n}\right)\right)\right)
\end{aligned}$$

187 **6 Parameters  $a$  and  $d$  in the example on page**  
188 **172 of Izhikevich (2007)**

189 The example on page 172 of Izhikevich (2007) calculates the *approximate* topo-  
190 logical normal form for the  $I_{Na_p} + I_K$  model

$$\begin{aligned}
\dot{V} &= \frac{1}{C}((I + I_0) - g_L(V + V_0 - E_L) \\
&\quad - g_{Na}m_\infty(V + V_0)(V + V_0 - E_{Na}) \\
&\quad - g_K(n + n_0)(V + V_0 - E_K)) = F(V, n, I) \quad (39) \\
\dot{n} &= \frac{n_\infty(V + V_0) - (n + n_0)}{\tau} = G(V, n, I) \quad (40)
\end{aligned}$$

<sup>191</sup> where the  $V$ ,  $n$ , and  $I$  have been centered at the Andronov-Hopf bifurcation  
<sup>192</sup> point  $V_0 = -56.4815$ ,  $n_0 = 0.0914$ , and  $I_0 = 14.659$ , respectively. To calcu-  
<sup>193</sup> late the value of the parameters  $a$  and  $d$  for the topological normal form in  
<sup>194</sup> Equations 6.8 and (6.9), we first compute the partial derivaties of  $F(V, n, I)$   
<sup>195</sup> and  $G(V, n, I)$  in Equations 39 and 40, respectively, evaluated at  $V = 0$ ,  $n = 0$ ,  
<sup>196</sup>  $I = 0$

$$\begin{aligned}
F_V &= \frac{1}{C}(-g_L - g_{Na}(m'_\infty(V_0)(V_0 - E_{Na}) + m_\infty(V_0)) - g_K n_0) \\
F_{VV} &= -\frac{g_{Na}}{C}(m''_\infty(V_0)(V_0 - E_{Na}) + m'_\infty(V_0)) \\
F_{VVV} &= -\frac{g_{Na}}{C}(m'''_\infty(V_0)(V_0 - E_{Na}) + 2m''_\infty(V_0)) \\
F_n &= \frac{1}{C}(-g_K(V_0 - E_K)) \\
F_{nn} &= 0 \\
F_{nnn} &= 0 \\
F_{Vn} &= -\frac{g_K}{C} \\
F_{VVn} &= 0 \\
F_{Vnn} &= 0 \\
G_V &= \frac{n'_\infty(V_0)}{\tau} \\
G_{VV} &= \frac{n''_\infty(V_0)}{\tau} \\
G_{VVV} &= \frac{n'''_\infty(V_0)}{\tau} \\
G_n &= -\frac{1}{\tau} \\
G_{nn} &= 0 \\
G_{nnn} &= 0 \\
G_{Vn} &= 0 \\
G_{VVn} &= 0 \\
G_{Vnn} &= 0
\end{aligned}$$

Next, we use these partial derivatives to obtain the partial derivatives in the previous section. Finally, we use the partial derivatives in the previous section to obtain the values of the parameters  $a$  and  $d$  from Equations 21 and 22, respectively. This gives  $a = -0.002970$  and  $d = -0.002613$ . Code for the calculation of the values of parameters  $a$  and  $d$  is given in:  
<https://sccn.ucsd.edu/~rapela/dynamicalSystems/izhikevich07/ch6/code/>

## 7 Evidence for typo in the example on page 172 of Izhikevich (2007)

In the example on page 172, Izhikevich (2007) reports parameters values  $a = -0.0026$  and  $d = -0.0029$  for the topological normal form in Equations 6.8 and (6.9) of the  $I_{Na,p} + I_K$  model. These parameters values are different from the ones derived in the previous section ( $a = -0.002970$  and  $d = -0.002613$ ). That the parameter value for  $a$  ( $d$ ) reported in Izhikevich (2007) is very similar to the parameter value for  $d$  ( $a$ ) derived in the previous section suggests a typo in Izhikevich (2007), where the values of parameters  $a$  and  $d$  have been interchanged. To check which set of parameters values is more adequate, the ones derived in the previous section or the ones provided in Izhikevich (2007), we compared descriptors (amplitude and frequency) of limit cycles in simulations of the  $I_{Na,p} + I_K$  model of the example on page 172 of Izhikevich (2007) with theoretical values of these descriptors computed from the parameters derived in the previous section and from the parameters given in Izhikevich (2007). We computed theoretical values using formulas on page 173 of Izhikevich (2007). Figures 1a and 1b are as the center and left panels, respectively, in Figure 6.12 in Izhikevich (2007). They show that the parameters values derived in the previous section yield theoretical values of the amplitude and frequency of limit cycles closer to descriptors extracted from model simulations than the parameter values provided in Izhikevich (2007). This indicates that the parameter values derived in the previous section are more adequate for the topological normal form in Equations 6.8 and 6.9 of the  $I_{Na,p} + I_K$  model of the example on page 172 of Izhikevich (2007) than later parameter values reported in Izhikevich (2007), specially for input currents closer to the bifurcation point,  $I_{ah} = 14.659$ .

We simulated the  $I_{Na,p} + I_K$  model in the example of page 172 of Izhikevich (2007) and the topological normal form in Equations 6.8 and 6.9 with the parameters derived in the previous section and with the parameters in Izhikevich (2007). For these simulations we used a constant injected dc-current  $I = 17$ . The blue trace in Figure 2 plots voltages from the  $I_{Na,p} + I_K$  model. The green and red traces plot voltages from the topological normal forms with the parameters derived in the previous section and with the parameters in Izhikevich (2007), respectively. We see that the voltages from the topological normal form with the parameters derived in the previous section better approximate the voltages

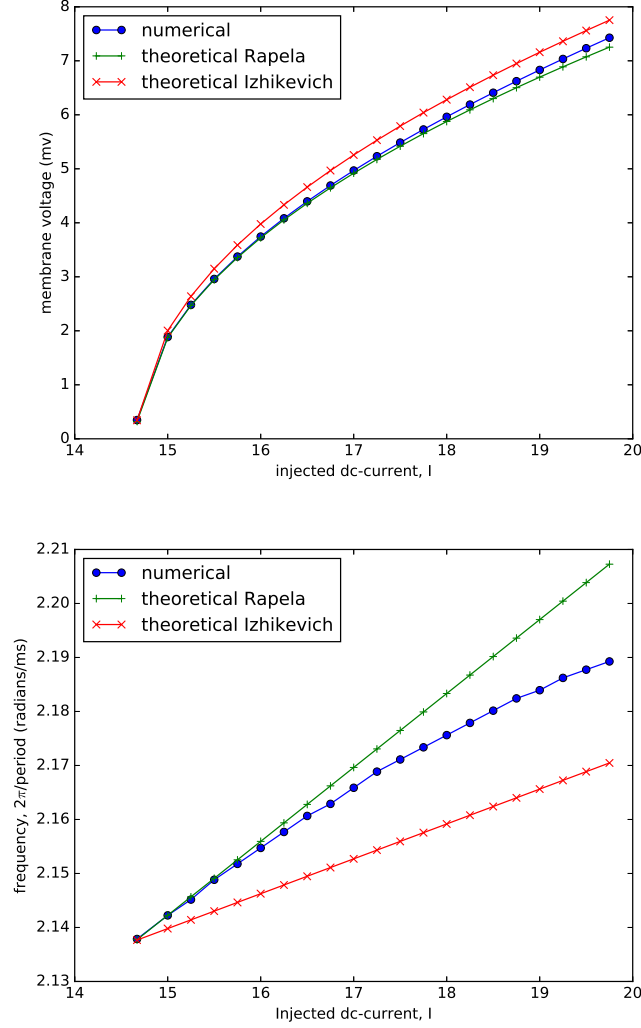


Figure 1: Descriptors (amplitude in panel (a) and frequency in panel (b)) of limit cycles next to an Andronov-Hopf bifurcation for the  $I_{Na,p} + I_K$  model on the example of page 172 of Izhikevich (2007) as a function of the input current to the model. Blue traces give these descriptors values derived from numerical simulations. Green traces provide theoretical values of these descriptors, from the formulas given on page 173 of Izhikevich (2007), using parameters  $a = -0.00297$  and  $d = -0.002613$  derived in the previous section. Red traces are as the green traces but for the parameters  $a = 0.0026$  and  $d = 0.0029$  used in the example of page 172 of Izhikevich (2007). Theoretical values obtained using the parameters  $a$  and  $d$  from the previous section better approximate descriptors derived from numerical simulations than theoretical values obtained using the parameters in Izhikevich (2007).

238 from the model than the voltages from the topological normal form with the pa-  
239 rameters in Izhikevich (2007). Thus, Figure 2 again shows that the parameters  
240  $a$  and  $d$  of the topological normal form in Equations 6.8 and 6.9 derived in the  
241 previous section are more adequate than the ones derived in Izhikevich (2007).

242

## 243 References

- 244 Izhikevich, E. M. (2007). *Dynamical systems in neuroscience*. MIT press.
- 245 Kuznetsov, Y. A. (2004). *Elements of applied bifurcation theory* (Vol. 112).  
246 Springer Science & Business Media.

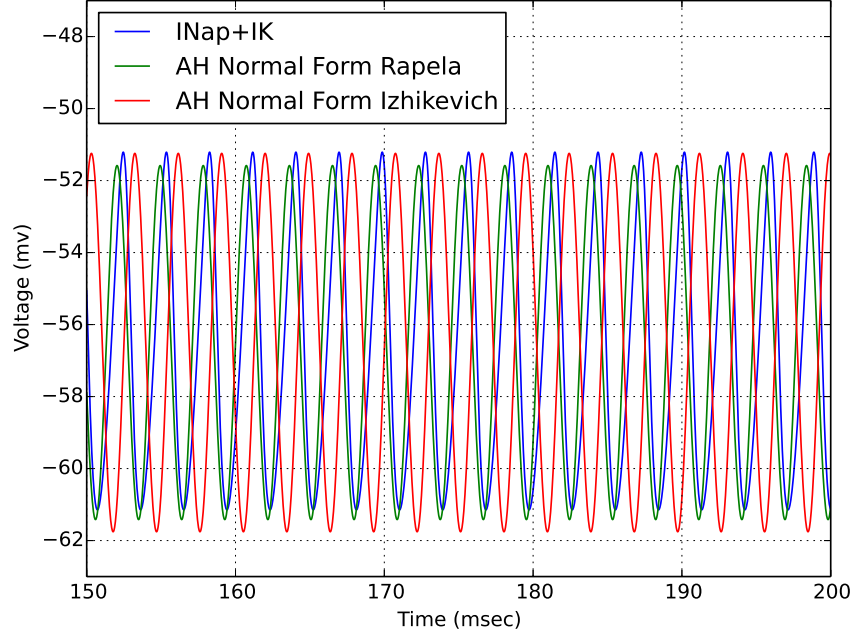


Figure 2: Voltages from the simulation close to an Andronov-Hopf bifurcation (input current  $I=17$ ) of the  $I_{Na,p} + I_K$  model on the example of page 172 of Izhikevich (2007) and from its topological normal form given in Equations 6.8 and 6.9 of Izhikevich (2007). The blue trace gives voltages from the simulation of the  $I_{Na,p} + I_K$  model. The green trace provides voltages from the simulation of the topological normal form using the parameters  $a = -0.00297$  and  $d = -0.002613$  derived in the previous section. The red trace is as the green one, but for the parameters  $a = 0.0026$  and  $d = 0.0029$  used in the example on page 172 of Izhikevich (2007). Voltages from the topological normal form with the parameters from the previous section are better approximations to the voltages from the simulation of the  $I_{Na,p} + I_K$  model than voltages from the topological normal form with the parameters used in the example on page 172 of Izhikevich (2007).